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ON THE NONEXISTENCE OF A SYMMETRY
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Abstract

It is shown that there no exist finite continuous groups of transformations in the Minkowsky space-time which contain the conformal group as a subgroup.

On the nonexistence of a symmetry

higher than conformal

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The Lorentz and Poincare groups play as known an important role as the groups of a symmetry of the quantum field theory. Lately there exists considerable interest also to the conformal group, which is an exact group of symmetry for massless particles and a broken group of symmetry for massive particles. The remarkable feature of the conformal symmetry is that it fixs the form of two— and threepoint functions independently of the concrete dynamics [1,2]. It would be attractive to find such group of transformations in the Minkowsky space—time which would define a form of the fourpoint functions and so on. Independently of the question of the fixing of a part of the dynamics by a group of the symmetry it is interesting to study possible space—time symmetries higher than Poincare and conformal.

The purpose of this paper is to show that there no exist finite continuous groups (Lie groups) of transformations in the Minkowsky space-time containing a subgroup which is locally isomorphic to the conformal group. Therefore, there no exists finite continuous group of space-time symmetry higher than conformal.

The generators of the continuous group of transformations in Minkowsky space-time have as known the form

$$L_{i}(x) = i \int_{i}^{i} \int$$

where the functions $f_{i\mu}(x)$ are determined from the transformation law of coordinates. In particular, for generators of the conformal group we have [3]:

$$M_{\mu\nu} = i(x_{\mu}\frac{\partial}{\partial x_{\nu}} - x_{\nu}\frac{\partial}{\partial x_{\mu}}), \quad P_{\mu} = i\frac{\partial}{\partial x_{\mu}},$$

$$K_{\mu} = i(x^{2}g_{\mu\nu} - 2x_{\mu}x_{\nu})\frac{\partial}{\partial x_{\nu}}, \quad \mathcal{D} = ix_{5}\frac{\partial}{\partial x_{5}}$$
(2)

We will show that in the Minkowsky space-time it isn't possible to realize the algebra of the finite continuous group which has the generators of form (1) and contain the conformal group as a subgroup.

It is convenient to use the isomorphism between the conformal group and the group SO(2,4). The generators Lab (a,b = 0,1,2,3,5,6) of the group SO(2,4) and generators of the conformal group are connected by the relations [3]:

$$L_{\mu\nu} = M_{\mu\nu}, L_{5\mu} = \frac{1}{2}(P_{\mu} + K_{\mu}),$$

$$L_{6\mu} = \frac{1}{2}(P_{\mu} - K_{\mu}), L_{56} = \mathcal{D}.$$
(3)

We consider, for visual aids, the particular case.

Let us assume that in the Minkowsky space-time there exists a group of transformations which is locally isomorphic to the group SO(P,9) (P+9>6, P>2, 9>4). with generators \mathcal{L}_{AB} (A,B=0,1,2,3,5,6,7,...,P+9), where $\mathcal{L}_{AB}=\mathcal{L}_{AB}$. Since

where $g_{AB} = (1,-1,-1,-1,1,...)$, among the generators \mathcal{L}_{AB} at least one generator Q (for example \mathcal{L}_{67}) which commute simultaneously with $\mathcal{L}_{\mu\nu}$ and $\mathcal{L}_{5\mu}$, which form algebra SO(1,4), always exist. Let us find a form of such generator. From the condition $[Q,\mathcal{L}_{\mu\nu}] = 0$ follows that

$$Q(x) = i \Phi(x^2) \times_{\nu} \frac{\partial}{\partial x^{\nu}}, \qquad (4)$$

where $\Phi(x^2)$ is arbitrary scalar function of $X = X_0^2 - \overline{X}^2$. Then from the condition $[Q, L_{5\mu}] = 0$ we find, taking into account (2-4):

$$(1-x^{2})g_{\mu\rho}\Phi(x^{2}) + 2x_{\mu}x_{\rho}[(1-x^{2})\Phi'(x^{2}) + \Phi(x^{2})] = 0,$$
(5)

$$\Phi'(z) = \frac{\partial \Phi(z)}{\partial z} ; \quad \mu, g = 0, 1, 2, 3$$

From the equations (5) we obtain $(1-x^2) \mathcal{P}(x^2) = 0$

and, therefore the solution of the equations (5) is $\phi = 0^{8}$.

Thus, the generator Q=0, i.e. in the Minkowsky space-time there no exists an operator of the form (1) which commute simultaneously with $M\mu\nu$ and $P\mu + K\mu$. Therefore, the group SO(P,q) (P+q>6, $P\geq 2$, $q\geq 4$) cannot be realize as a group of coordinate transformations in the Minkowsky space-time.

The proof listed above is generelazed to arbitrary finite continuous group which contain the conformal subgroup. Indeed, any finite Lie group is locally isomorphic to some linear group [4]. And for linear groups as it isn't difficult to show at least one generator which does not belong to the subalgebra $SO(2,4) \sim SU(2,4)$ and commute with generators $L_{\mu\nu}$ and $L_{5\mu}$ of the subgroup $SO(4,4) \sim USP(2,2)$ (where $SO(4,4) \subset SO(2,4)$) always exists. But as we have seen such generator cannot be realized in the Minkowsky space-time.

Thus, the finite continuous group which contain the conformal group as a subgroup cannot be realized as the group of coordinate transformations in the Minkowsky spacetime. Therefore, there no exists finite continuous group

of space-time symmetry higher than conformal *.

There exist at least one infinite continuous group, which contain the conformal group as a subgroup - it is the group of general coordinate transformations, recently considered in [5]. We would like to note, however, that the structure of general coordinate transformations group is essentially distinguished from the structure of Lie groups (for example, SO(P, 9), SU(P, 9)) in the formal limit of infinite number of parameters $(P, 9 \rightarrow \infty)$.

We see, thus, that finite group of symmetry can fix only small part of the dynamics, maximally two- and three-point functions (conformal group). Therefore, in order to fix the more part of the dynamics we must consider an infinite groups.

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We exclude singular solution $\varphi(x^2)$ =

 $^{= \}delta(1-x^2) \, \varphi(x^2) \qquad , \text{ where } \delta(z) \text{ is a delta-}$ function, $\varphi(x^2)$ is arbitrary function.

For space with a metric which contain n plus and m minus the maximal group of space symmetry compatible with Lorentz invariance is the conformal group of this space, i.e. the group SO(n+1,m+1).

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